

# Number-Rigidity and $\beta$ -Circular Riesz Gas

**David Dereudre**

(joint work with Thibaut Vasseur)

Laboratoire de Mathématiques Paul Painlevé,  
University of Lille, France

Conference "Coulomb gases and universality",  
December 2022, Paris

# Introduction

- Gibbs point process on  $\mathbb{R}^d$  interacting with the Riesz pair potential

$$g(x) = \frac{1}{|x|^s} \quad d-1 < s < d$$

- $g$  is non-integrable at infinity,  $\nabla g$  is integrable.
- canonical ensemble with constant density  $\rho > 0$  and inverse temperature  $\beta > 0$ .
- periodic boundary condition
- number-rigidity and equivalence of ensembles

- 1 The Model
- 2 Number-Rigidity
- 3 Equivalence of ensembles

# 1 The Model

# The Riesz energy with background

$\gamma = \{x_1, \dots, x_n\}$  included  $\Lambda_n = [-n^{1/d}/2, n^{1/d}/2]^d$

$$H(\gamma) = \sum_{\{x,y\} \in \gamma} g(x-y) = \frac{1}{2} \int \int_{\mathbb{R}^d \setminus \text{Diag}} g(x-y) \gamma(dx) \gamma(dy).$$

With the background on  $\Lambda_n$

$$\tilde{H}_n(\gamma) = \frac{1}{2} \int \int_{\Lambda_n \setminus \text{Diag}} g(x-y) (\gamma(dx) - dx) (\gamma(dy) - dy).$$

The energy  $\tilde{H}_n(\gamma)$  is of order  $n$  (the volume).

# The periodic Riesz energy

For  $k \geq 1$ ,  $\gamma^k$  is the concatenation of  $(2k+1)^d$  copies of  $\gamma$  in the translations of  $\Lambda_n$ . It is a configuration in  $\Lambda_{(2k+1)^d n}$ .

## Proposition

$$\lim_{k \rightarrow \infty} \frac{\tilde{H}_{\Lambda_{(2k+1)^d n}}(\gamma^k)}{(2k+1)^d} = \sum_{\{x,y\} \in \gamma} g_n(x-y) + n\varepsilon_n$$

with  $g_n(x) = \sum_{k \in \mathbb{Z}^d} (g(x + kn^{1/d}) - \frac{1}{n} \int_{\Lambda_n} g(y + kn^{1/d}) dy)$ .

For all  $x \in \Lambda_n$ ,  $|g_n(x) - g(x)| \leq Cn^{-s/d}$ .

## Definition

The periodic Riesz energy of  $\gamma$  in  $\Lambda_n$  is defined by

$$H_n(\gamma) = \sum_{\{x,y\} \subset \gamma} g_n(x-y).$$

Properties of  $g_n$ 

$$g_n(x) = \sum_{k \in \mathbb{Z}^d} (g(x + kn^{1/d}) - \frac{1}{n} \int_{\Lambda_n} g(y + kn^{1/d}) dy).$$

## Proposition

- (Stability) There exists a constant  $A \geq 0$  such that for point configuration  $\gamma \in \Lambda_n$  such that  $|\gamma| = n$ , we have  $H_n(\gamma) \geq -An$ .
- (Shift invariance) For every  $u \in \Lambda_n$  and every configuration  $\gamma$  in  $\Lambda_n$  we have  $H_n(\tau_u^n(\gamma)) = H_n(\gamma)$ .
- (Approximation) There exists a constant  $c > 0$  such that for every point  $x \in \Lambda_n$  we have

$$|g_n(x) - g(x)| \leq cn^{-s/d}$$

# The canonical ensemble

$\text{Bin}_{\Lambda,n}$  is the distribution of  $n$  independent points uniformly distributed in  $\Lambda$ .

## Definition

The canonical Gibbs measure in  $\Lambda_n$  with inverse temperature  $\beta > 0$  is

$$P_n^\beta = \frac{1}{Z_n^\beta} e^{-\beta H_n} \text{Bin}_{\Lambda_n, n}.$$

## Theorem

*The sequence  $(P_n^\beta)_{n \geq 1}$  admits accumulation points for the local convergence topology. They are called  $\beta$ -circular Riesz gases.*

Uniqueness or non-uniqueness of accumulation points is unknown (excepted for  $d = 1$ , Boursier 2022).



# Arguments for tightness

- The energy is stable : For any  $\gamma$  such that  $\#\gamma = n$

$$H_n(\gamma) \geq -An.$$

- The partition function : there exists  $0 < a_\beta < b_\beta < +\infty$

$$a_\beta^n \leq Z_n^\beta \leq b_\beta^n.$$

- The relative entropy is uniformly bounded

$$I(P_n^\beta | \pi_{\Lambda_n}) / |\Lambda_n| \leq C.$$

- $P_n^\beta$  is stationary on the torus  $\Lambda_n$ .

# Some references

- Gruber, Lugin, Lieb, Lebovitz, Martin (1975-85) :  
Thermodynamic limit, charge neutrality, screening, sum rules.
- Hardin, Saff and Simanek (2014) : Periodic energy of a crystal
- Leblé-Serfaty (2017) : LDP for Riesz gases with confining potential
- Valko, Virag (2009), Killip-Stoiciu (2009), Nakano (2014) :  
 $\beta$ -circular ensembles and the Sine- $\beta$  process ( $s = 0$ ,  $d = 1$ )
- Boursier (2022) : Riesz gas on the circle ( $0 < s < 1$ ,  $d = 1$ )
- Lewin (2022) : Survey on Riesz gas.

## 2 Number-Rigidity

# Number-Rigidity

## Definition (Ghosh-Peres 2017)

A point process  $\Gamma$  in  $\mathbb{R}^d$  is said number-rigid if for any bounded set  $\Lambda \subset \mathbb{R}^d$  there exists a function  $F_\Lambda$  such that almost surely

$$\#\Gamma_\Lambda = F_\Lambda(\Gamma_{\Lambda^c}).$$

Are the  $\beta$ -circular Riesz gases number-rigid?

Previous works for Gibbs point process :

- $s > d$  summable potential : Non number-rigidity (grand canonical DLR equations)
- $s = 0, d = 2$  and  $\beta = 2$  : Number-Rigidity (DPP structure + linear statistics), Ghosh-Lebowitz 2017
- $s = 0, d = 1$  and  $\beta > 0$  : Number-Rigidity (canonical DLR equations or linear statistics), D.-Leblé-Hardy-Maïda 2019 or Chhaibi-Najnudel 2018.

# One point deletion

## Definition (Holroyd-Soo 2013)

A point process  $\Gamma$  in  $\mathbb{R}^d$  is said one-point deletion if for any random variate  $X \subset \Gamma$  the distribution of  $\Gamma \setminus X$  is absolutely continuous with respect to  $\Gamma$ .

"Non number-rigidity" and "One point deletion" are almost equivalent.

Heuristically, for Gibbs point processes and if  $X$  is "typical"

$$\frac{P_{\Gamma}}{P_{\Gamma \setminus X}} = C^{st} e^{-\beta h(X, \Gamma \setminus X)}.$$

The one point deletion property requires a good definition for

$$h(X, \Gamma \setminus X).$$

# The energy of a point

Let  $x \in \mathbb{R}^d$  and  $\gamma$  an infinite configuration ( $x \notin \gamma$ ). Three candidates for  $h(x, \gamma)$  :

$$h_1(x, \gamma) = \sum_{y \in \gamma} \frac{1}{|x - y|^s} = \int \frac{1}{|x - y|^s} \gamma(dy)$$

$$h_2(x, \gamma) = \lim_{n \rightarrow \infty} \int_{\Lambda_n} \frac{1}{|x - y|^s} (\gamma(dy) - dy)$$

$$h_3(x, \gamma) = \lim_{n \rightarrow \infty} \left( \int_{\Lambda_n} \frac{1}{|x - y|^s} \gamma(dy) - C_n(\#\gamma_{\Lambda_n}, \gamma_{\Lambda_n^c}) \right)$$

# The main result

## Theorem (D. and Vasseur 2022)

*For any  $\beta > 0$ , there exists a  $\beta$ -circular Riesz gas  $P_\star^\beta$  which is not number-rigid.  $P_\star^\beta$  is also one-point deletion.*

$P_\star^\beta = \lim_{k \rightarrow \infty} P_{n_k}^\beta$  for a subsequence  $(n_k)$ .

We believe that all  $\beta$ -circular Riesz gases are not number-rigid.

## Corollary

*For any bounded  $\Lambda$  and  $k \geq 0$  then for all  $P_\star^\beta$ -a.s.  $\gamma$ ,*

$$P_\star^\beta(N_\Lambda = k | \gamma_{\Lambda^c}) > 0.$$

## Main ingredient of the proof

## Proposition

For any  $\beta > 0$ , there exists a constant  $K > 0$  and a subsequence  $(n_k)_{k \geq 1}$  such that for all  $k \geq 1$

$$\int |h_{n_k}(0, \gamma)| P_{n_k}^\beta(d\gamma) \leq K,$$

where

$$h_n(x, \gamma) = \sum_{y \in \gamma} g_n(x - y),$$

$$g_n(x) = \sum_{k \in \mathbb{Z}^d} (g(x + kn^{1/d}) - \frac{1}{n} \int_{\Lambda_n} g(y + kn^{1/d}) dy).$$



### 3 Equivalence of ensembles

## General principle

**Canonical ensembles :** The density of particles  $\rho > 0$  is prescribed. During the thermodynamic limit ( $\Lambda_n \rightarrow \infty$ ) the number of particles is fixed equal to  $\rho|\Lambda_n|$ .

**Grand canonical ensembles :** The activity  $z > 0$  (or the chemical potential  $\mu$ ) is prescribed ( $z = e^{-\beta\mu}$ ). During the thermodynamic limit ( $\Lambda_n \rightarrow \infty$ ) the number of particles is random. The Gibbs process is absolutely continuous with respect to the Poisson point process with intensity  $z > 0$ .

### Definition (Equivalence of ensembles)

The canonical ensembles and the grand canonical ensembles are the same. There exist functions  $\rho \mapsto z_\rho$  and  $z \mapsto \rho_z$ .

The equivalence of ensembles is proved for a large class of summable pairwise potentials (Ruelle 70, Georgii 94, Vasseur 2022), including the Riesz potential for  $s > d$ .

# Equivalence of ensembles with the DLR formalism

- A canonical ensemble  $P$  satisfies the canonical DLR (Dobrushin-Lanford-Ruelle) equations :

$$P(d\gamma_\Lambda | \#\gamma_\Lambda = k, \gamma_{\Lambda^c}) = \frac{1}{Z_\Lambda^\beta(k, \gamma_{\Lambda^c})} e^{-\beta H(\gamma_\Lambda | \gamma_{\Lambda^c})} \text{Bin}_{\Lambda, k}(d\gamma_\Lambda).$$

- A grand canonical ensemble  $P$  satisfies the grand canonical DLR equations :

$$P(d\gamma_\Lambda | \gamma_{\Lambda^c}) = \frac{1}{Z_\Lambda^\beta(\gamma_{\Lambda^c})} e^{-\beta H(\gamma_\Lambda | \gamma_{\Lambda^c})} \pi_\Lambda^z(d\gamma_\Lambda).$$

## Definition (Equivalence of ensembles)

If  $P$  satisfies the canonical DLR equations then  $P$  satisfies the grand canonical DLR equations.

# Canonical DLR equations for $\beta$ -circular Riesz gas

The energy to move a particle from 0 to  $x$  in  $\gamma$  is

$$M(x|\gamma) = \sum_{y \in \gamma} g(x - y) - g(y).$$

## Theorem (Canonical DLR equations)

Let  $\mathbb{P}^\beta$  be a  $\beta$ -Circular Riesz gas,  $\Lambda$  be a bounded Borel subset of  $\mathbb{R}^d$ . Then for  $\mathbb{P}^\beta$  a.e.  $\gamma$

$$P^\beta(d\gamma_\Lambda | \#\gamma_\Lambda = k, \gamma_{\Lambda^c}) = \frac{1}{Z_\Lambda^\beta(k, \gamma_{\Lambda^c})} e^{-\beta H(\gamma_\Lambda | \gamma_{\Lambda^c})} \text{Bin}_{\Lambda, k}(d\gamma_\Lambda).$$

where  $H(\gamma_\Lambda | \gamma_{\Lambda^c}) = \sum_{\{x, y\} \subset \gamma_\Lambda} g(x - y) + \sum_{x \in \gamma_\Lambda} M(x | \gamma_{\Lambda^c})$ .

Similar proof as D., Leblé, Hardy and Maïda (2019) for the Sine- $\beta$  process.

# Grand canonical DLR equations for $P_\star^\beta$

Based on the one-point deletion property of  $P_\star^\beta$

## Theorem (Grand canonical DLR equations)

Let  $\Lambda$  be a bounded Borel subset of  $\mathbb{R}^d$ . Then for  $P_\star^\beta$  a.e.  $\gamma$

$$P_\star^\beta(d\gamma_\Lambda | \gamma_{\Lambda^c}) = \frac{1}{Z_\Lambda^\beta(\gamma_{\Lambda^c})} e^{-\beta H(\gamma_\Lambda | \gamma_{\Lambda^c})} \pi_\Lambda(d\gamma_\Lambda).$$

where  $\gamma_\Lambda = \{x_1, x_2, \dots, x_k\}$  and

$$H(\gamma_\Lambda | \gamma_{\Lambda^c}) = h(x_1, \gamma_{\Lambda^c}) + h(x_2, x_1 \cup \gamma_{\Lambda^c}) + \dots + h(x_k, x_1 \cup \dots \cup x_{k-1} \cup \gamma_{\Lambda^c})$$

$$h(x, \gamma) = \lim_{n \rightarrow \infty} \left( \sum_{y \in \gamma_{\Lambda_n}} g(x - y) - C_n(\#\gamma_{\Lambda_n}, \gamma_{\Lambda_n^c}) \right).$$

## Integral compensator

$$h(x, \gamma) = \lim_{n \rightarrow \infty} \left( \sum_{y \in \gamma_{\Lambda_n}} g(x - y) - C_n(\#\gamma_{\Lambda_n}, \gamma_{\Lambda_n^c}) \right).$$

We believe that the integral compensator works

$$C_n(\#\gamma_{\Lambda_n}, \gamma_{\Lambda_n^c}) = \int_{\Lambda_n} g(y) dy$$

## Proposition

*If  $P^\beta$  is hyperuniform with  $\text{Var}(N_\Lambda) \leq C|\Lambda|^{s/d-\varepsilon}$  then the grand canonical DLR equations hold with*

$$h(x, \gamma) = \lim_{n \rightarrow \infty} \left( \sum_{y \in \gamma_{\Lambda_n}} g(x - y) - \int_{\Lambda_n} g(y) dy \right).$$

It is the case for  $d = 1$ , Boursier 2022.

# Summary of the talk

- We define infinite volume Riesz gases ( $d - 1 < s < d$ ) in  $\mathbb{R}^d$  at inverse  $\beta > 0$  with periodic boundary conditions.
- At least one of them  $P_\star^\beta$  is not number Rigid.
- $P_\star^\beta$  satisfies canonical and grand canonical DLR equations.
- The energy of a point  $x$  in  $\gamma$  exists

$$h(x, \gamma) = \lim_{n \rightarrow \infty} \left( \sum_{y \in \gamma_{\Lambda_n}} g(x - y) - C_n(\#\gamma_{\Lambda_n}, \gamma_{\Lambda_n^c}) \right).$$

- If  $d = 1$ ,  $P_\star^\beta$  is hyperuniform (Boursier 2022) and so

$$h(x, \gamma) = \lim_{n \rightarrow \infty} \left( \sum_{y \in \gamma_{\Lambda_n}} g(x - y) - \int_{\Lambda_n} g(y) dy \right).$$

# Open questions

- Hyperuniformity and integral compensator for  $d \geq 2$ .
- DLR equations for  $s \leq d - 1$ ?
- Does the Number-Rigidity property appear at  $s = 0$ ? (true for  $d = 1$ )