

Some aspects of the Anderson hamiltonian with white noise in 1D

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Coulomb gases and universality

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Motivations

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Ultimate goal : Understand random Schrödinger operators using the link with RMT ?

Schrödinger operator

Differential self-adjoint operator on \mathbb{R}^d :

$$u \mapsto -\Delta u + V \cdot u.$$

V : potential.

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V : potential.

Interpolation between Laplacian:

$$u \mapsto -\Delta u,$$

and multiplication by the potential V :

$$u \mapsto V \cdot u.$$

Schrödinger operator

Laplacian operator:

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electron conduction in a crystal.

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Main results/conjectures:

- ▶ $d = 1$: Anderson localization (result).
- ▶ $d = 2$: Anderson localization (conjecture).
- ▶ $d = 3$ and above, there is a delocalized phase if V small enough (major conjecture!).

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- ▶ **Advantages:** white noise is natural, it is explicit, and it has nice properties.
- ▶ **Inconvenient:** it is irregular.

Introduce:

$$\mathcal{H}_L : u \in \mathcal{D}_L \mapsto -\partial_x^2 u + \xi \cdot u.$$

where \mathcal{D}_L subdomain of $L^2(-L/2, L/2)$ with Dirichlet boundary conditions.

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And

$$\mathcal{H} : u \in \mathcal{D} \mapsto -\partial_t^2 u + \xi \cdot u.$$

where \mathcal{D} subdomain of $L^2(\mathbb{R})$.

Anderson localization

Theorem (L.D., C. Labbé '22+)

When $L \rightarrow \infty$, \mathcal{H}_L converges towards \mathcal{H} in the strong resolvent sense.

\mathcal{H} has a pure point spectrum, its eigenfunctions decay exponentially: for every eigenvalue λ ,

$$\frac{1}{|t|} \ln \sqrt{\varphi_\lambda(t)^2 + \varphi'_\lambda(t)^2} \xrightarrow{|t| \rightarrow \infty} -\gamma_\lambda$$

where $\gamma_\lambda > 0$ is the Lyapunov exponent associated to the eigenproblem.

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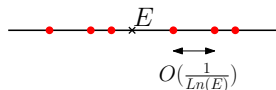
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This result was expected but there are some technical issues as white noise is singular.

Study of \mathcal{H}_L

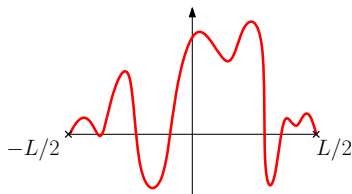
Study the fine statistics of the **spectrum** of \mathcal{H}_L when $L \rightarrow \infty$.

Eigenvalues



Random point process on the real line

Eigenvectors



Random real function on $[-L/2, L/2]$

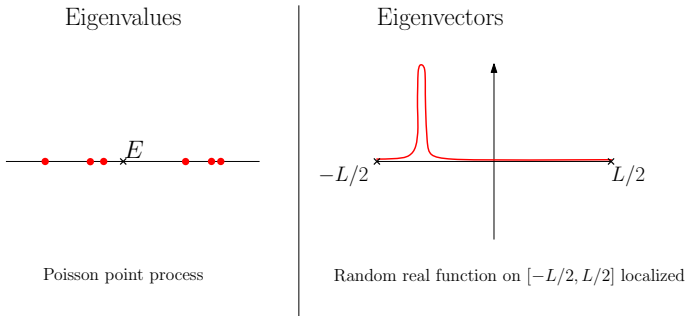
E : energy: it may depend on the size of the system L .

$n(E)$: density of states

(roughly, # of eigenvalues $[E, E + \varepsilon] \simeq n(E)L\varepsilon$)

Goal

Study the **spectrum** of this operator when $L \rightarrow \infty$.

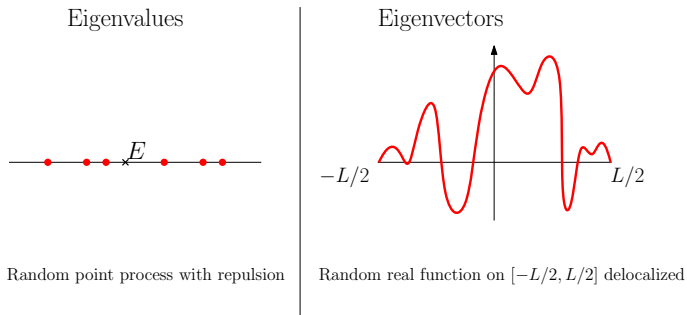


Usually for random operators with discrete spectrum, there is a dichotomy:

► **Localization** of the eigenvectors and **Poisson distribution** of eigenvalues.

Goal

Study the **spectrum** of this operator when $L \rightarrow \infty$.



Usually for random operators with discrete spectrum, there is a dichotomy:

► **Delocalization** of the eigenvectors and **repulsion** of the eigenvalues.

Summary of our results

- ▶ **When $E \ll L$:**

Eigenvalues: Poisson, *Eigenvectors:* localized.

- ▶ **When $E = O(L)$ or $E \gg L$:**

Eigenvalues: repulsion, *Eigenvectors:* delocalized.

Critical Regime $E = O(L)$

Zoom around energies E of order L .

► **Eigenvalues** :

$$Ln(E)(\lambda_i - E),$$

Critical Regime $E = O(L)$

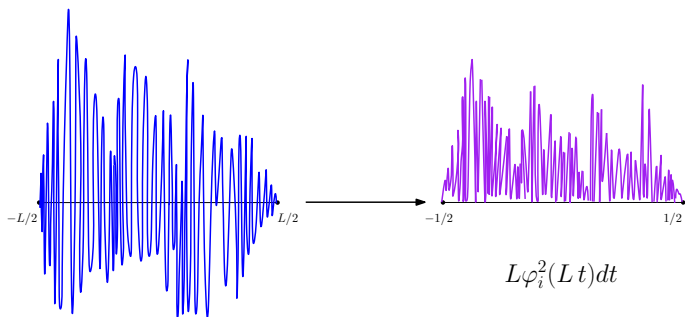
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► **Eigenvalues :**

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► **Eigenfunctions :**

$$L\varphi_i^2(Lt)dt.$$



Eigenfunction φ_i

Delocalization: Critical regime

Fix $\tau > 0$ and let $L/E \rightarrow \tau$.

$$Q_L := \sum_{i \geq 1} \delta_{(\operatorname{Ln}(E)(\lambda_i - E), L\varphi_i^2(Lt)dt)},$$

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\mathcal{Q}_L converges towards a random point process, characterized by coupled diffusions.

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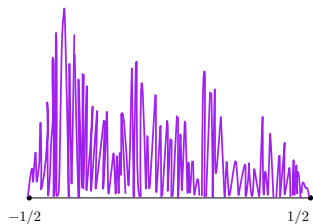
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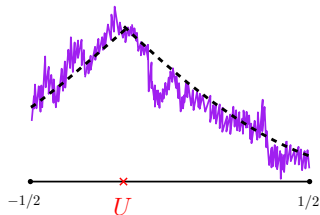
- ▶ Limiting eigenvalue point process = Sch_τ point process introduced by [Kritchevski, Valkó and Virág \(2012\)](#), in the context of discrete random Schrödinger operators. **Repulsion** between the points!
- ▶ Typical limiting eigenvectors = exponential of a Brownian motion plus a drift around a uniform point in $[-1/2, 1/2]$ (conjectured for this model by [Rifkind and Virág \(2018\)](#)). **Interaction** between the eigenvectors.

Delocalization: Behavior of eigenvectors



$$L\varphi_i^2(Lt)dt$$

Eigenvector associated to
an energy near E

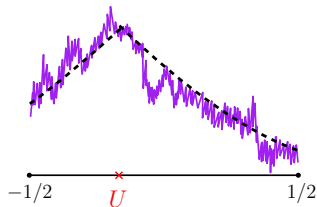


$$C \exp\left(-\frac{\tau}{4}|t - U| + \frac{\sqrt{\tau}}{\sqrt{2}}W(t - U)\right)dt$$

$U \sim$ uniform

W two-sided Brownian motion

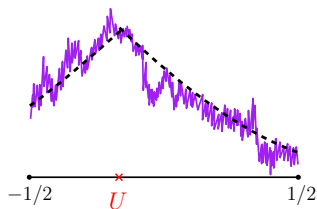
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Universal shape: conjectured by [Rifkind and Virág](#) to arise in many 1D models through their transition from localized to delocalized.

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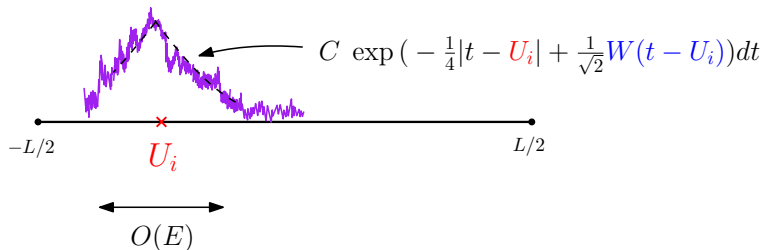
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This shape starts to appear in the **localized regime** when $1 \ll E \ll L$ when we zoom around the argmax of the function!

From localization to delocalization

Localized regime for $1 \ll E \ll L$: zoom around the argmax of the function!



Delocalization: Critical regime

Convergence at the level of **operators**:

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Theorem (D., Labbé ('21))

The rescaled (around E) and *unitarily changed operator* \mathcal{H}_L converges towards the operator

$$\text{CS}_\tau := 2 \begin{pmatrix} 0 & -\partial_t \\ \partial_t & 0 \end{pmatrix} + \sqrt{\tau} \begin{pmatrix} d\mathcal{B} + \frac{1}{\sqrt{2}}d\mathcal{W}_1 & \frac{1}{\sqrt{2}}d\mathcal{W}_2 \\ \frac{1}{\sqrt{2}}d\mathcal{W}_2 & d\mathcal{B} - \frac{1}{\sqrt{2}}d\mathcal{W}_1 \end{pmatrix},$$

where $\mathcal{B}, \mathcal{W}_1, \mathcal{W}_2$ are independent Brownian motions, on $L^2((0, 1), \mathbb{R}^2)$ with Dirichlet b.c..

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- ▶ The convergence holds in law in the strong resolvent sense + there is convergence in law of the eigenvalues/eigenvectors.
- ▶ Such a form was appearing in the works of [Edelman and Sutton \(2006\)](#) as a conjecture for the bulk of limiting random matrices.

Delocalization: Critical regime

Some remarks:

- ▶ The limiting operator acts on \mathbb{R}^2 valued functions: The initial space is **enlarged**. \rightarrow transformation of \mathcal{H}_L is non trivial.

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- ▶ The limiting operator acts on \mathbb{R}^2 valued functions: The initial space is **enlarged**. \rightarrow transformation of \mathcal{H}_L is non trivial.
- ▶ CS_τ is not properly defined in this form. An elegant way to define it is through *Dirac operators*, which appeared recently in [Valkó and Virág \(2016\)](#) for the limit of many random matrix models.
- ▶ Our approach could probably be extended to other models (bulk of tridiagonal matrices or other differential operators).

Link with the operators of Valkó and Virág

Dirac operators: differential self-adjoint operator of the form

$$\tau_R : f = \begin{pmatrix} f_1 \\ f_2 \end{pmatrix} \mapsto 2 R(t)^{-1} J \partial_t \begin{pmatrix} f_1 \\ f_2 \end{pmatrix}$$

$R(t)$: 2×2 symmetric positive matrix of $\det = 1$, $J = \begin{pmatrix} 0 & -1 \\ 1 & 0 \end{pmatrix}$.

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Valkó and Virág ('16) proved that the spectrum of those operators for a well-chosen random R corresponds to the spectrum of various famous RMT (Sine_β , Bessel_β etc.).

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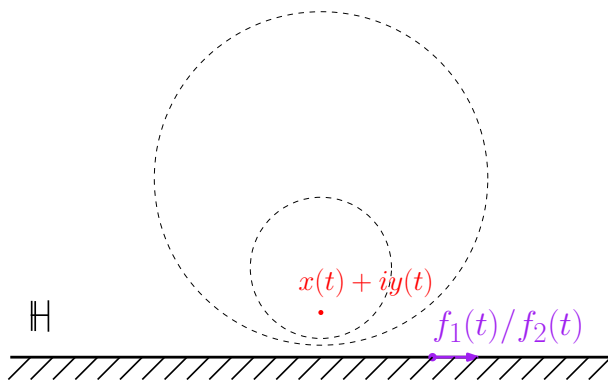
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Fact: $f_1(t)/f_2(t) \in \mathbb{R} \cup \{\infty\} = \partial\mathbb{H}$ follows hyperbolic rotation at speed λdt around $x(t) + iy(t)$ in \mathbb{H} .

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Link with the operators of Valkó and Virág

Theorem (D.,Labbé, 2021)

The operator CS_τ is unitarily equivalent to a Dirac operator associated to a hyperbolic Brownian motion of variance $\tau/2$.

Unitary map is random (non trivial).

Boundary conditions are complicated but explicit.

SOME IDEAS FOR THE PROOFS IN THE
LOCALIZED REGIME

Eigenvalue equation

Eigenvalue equation for \mathcal{H}_L defined on $[-L/2, L/2]$:

$$-\varphi'' + \xi \cdot \varphi = \lambda \varphi$$

with $\varphi(-L/2) = 0$ (without any condition on $\varphi(L/2)$).

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To keep in mind: it is much easier to analyse the solution of the ODE than the solution of the eigenvalue problem: λ is random and depends on the whole potential ξ .

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One can also impose first $\hat{\varphi}(L/2) = 0$ and solve

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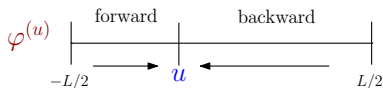
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Concatenation forward/backward

Key idea: Use *forward solution* φ_λ on the time-interval $[-L/2, u]$ and then *backward solution* $\hat{\varphi}_\lambda$ on $[u, L/2]$ for some well-chosen u .

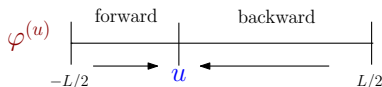
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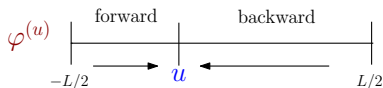


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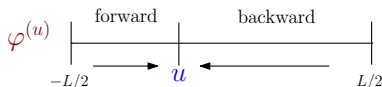
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It helps A LOT because it is **much easier** to analyze the **forward or backward solution of the ODE** than the eigenvalue equation (when λ eigenvalue, λ is random and depends on the whole potential $\xi!$).

Localization: A key formula

Proposition (Goldsheid Molchanov Pastur formula)

For all continuous, compactly supported on the first variable, and bounded G :

$$\begin{aligned} & \mathbb{E} \left[\sum_{\lambda \text{ eigenvalue}} G(\lambda, \varphi_\lambda) \right] \\ &= \int_{\lambda \in \mathbb{R}} \int_{-L/2}^{L/2} \int_0^\pi \sin^2(\theta) p_\lambda(\theta) p_\lambda(\pi - \theta) \mathbb{E} \left[G \left(\lambda, \frac{\varphi_\lambda^{(u)}}{\|\varphi_\lambda^{(u)}\|_2} \right) \right] d\lambda du d\theta, \end{aligned}$$

where

- ▶ $\varphi^{(u)}$ is the concatenation of the forward process and backward process at time u .
- ▶ $p_\lambda(\theta)$ transition probability of θ_λ “phase function” (argument of $\varphi'_\lambda + i\varphi_\lambda$).

THANK YOU!

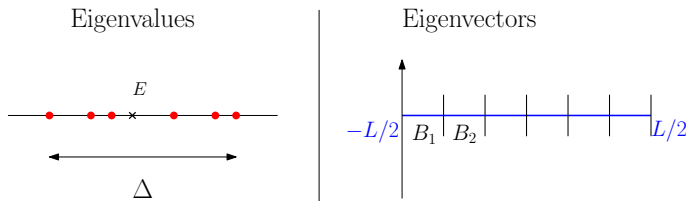
Strategy to prove convergence towards a Poisson point process when you know localization

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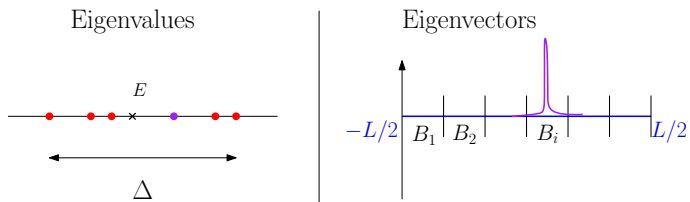
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Divide $[-L/2, L/2]$ into small boxes B_i , $i = 1, \dots, N$ of same length.

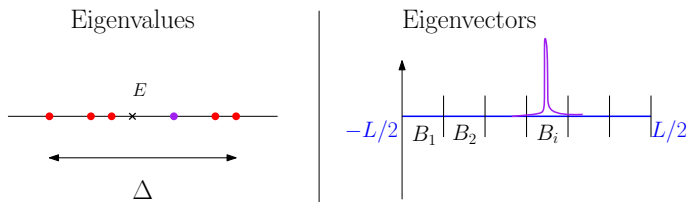


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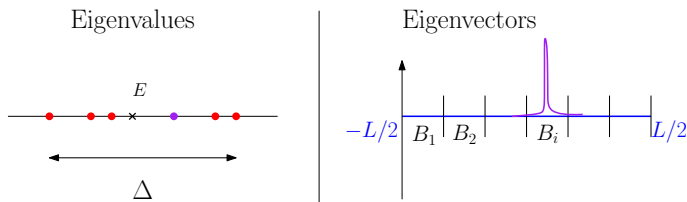


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- (C) $\sum_i \mathbb{P}[N_i(\Delta) \geq 2] \rightarrow 0$.